



Attribution & Literature

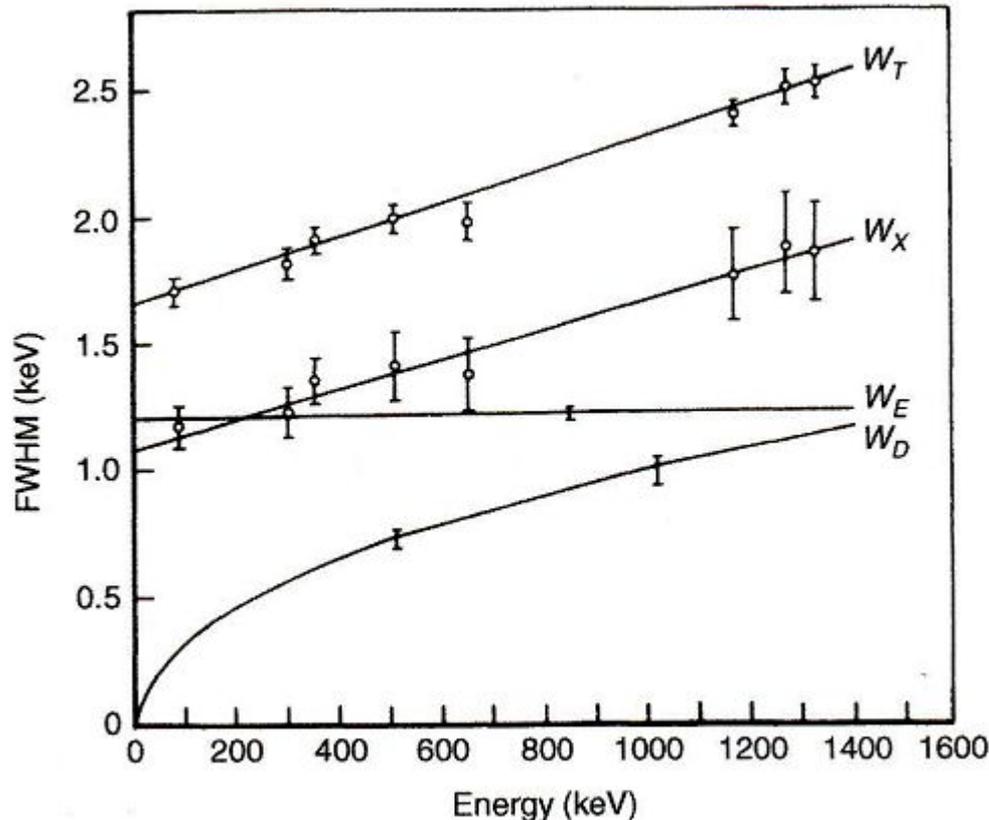
The material for this lecture is derived from [Spieler](#), particularly sections II, III, and IV. For the most complete treatment, consult Spieler's textbook [Semiconductor Detector Systems](#)



Energy Resolution

$$(\Delta E_{\text{total}})^2 = (\Delta E_{\text{stat}})^2 + (\Delta E_{\text{electronics}})^2 + (\Delta E_{\text{charge-loss}})^2$$

N.B. - More generally, shape of peak given by convolution of distributions describing each of these components!



Knoll 12.11

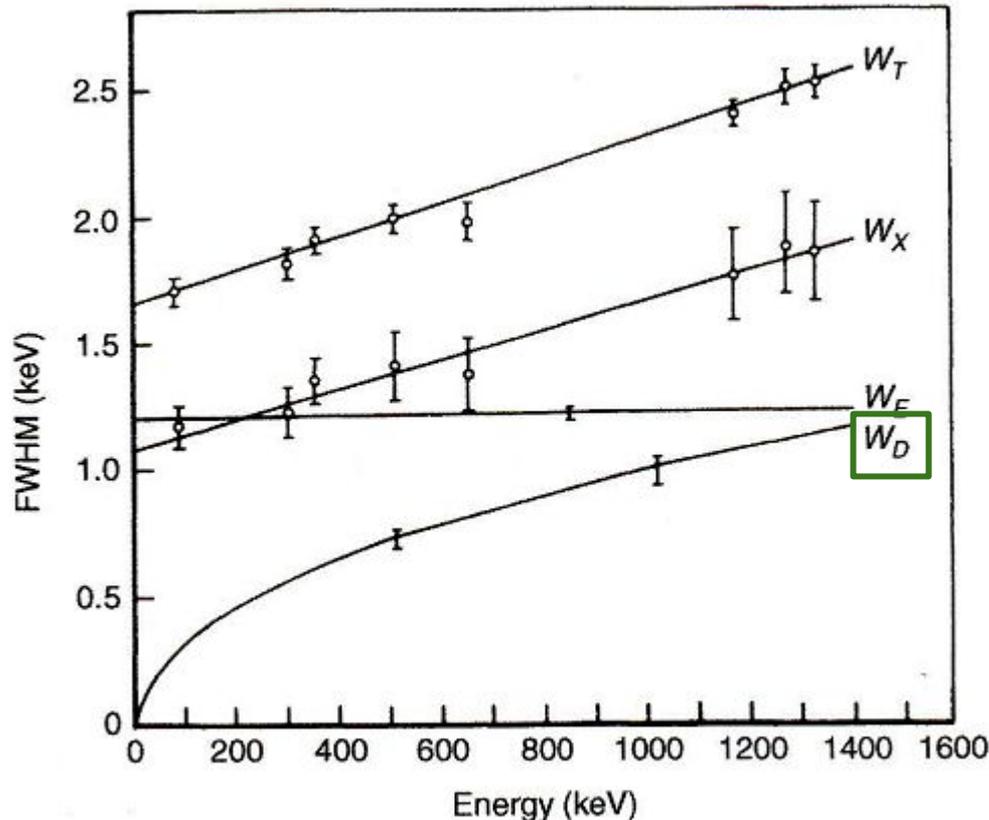
- Multiple elements contribute to width of peaks in energy spectrum
 - Statistics of carrier generation
 - Electronic Noise
 - Charge collection

Energy Resolution - Statistics of Carrier Generation



$$(\Delta E_{\text{total}})^2 = (\Delta E_{\text{stat}})^2 + (\Delta E_{\text{electronics}})^2 + (\Delta E_{\text{charge-loss}})^2$$

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Knoll 12.11

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Statistics of Carrier Generation

- Ionization spectrometers based on collection of N carriers generated by energy deposition E
- Carrier generation is a **stochastic process**
 - Let ε represent average energy required to generate an “information carrier”, IC
 - **Semiconductors:** IC = e/hole pairs | $\varepsilon = 3\text{-}5$ eV
 - **Gas Ionization:** IC = e/ion pairs | $\varepsilon \sim 30$ eV
 - **Scintillation:** IC = optical photons | $\varepsilon^* \sim 10$ eV
 - When accounting for collection/photocathode efficiency, $\varepsilon \sim 100$ eV
- The generation of charge carriers is often modelled with the **Poisson Distribution**



Statistics of Carrier Generation

- **Poisson model** has nice property where $\text{var}(N) = \sigma_N^2 = N$
- However, model only applicable if carrier generation is **independent** of all other carrier generation events
 - Good model for scintillators: many competing decay modes
 - More detail in scintillator lectures
 - For gas ionization and semiconductors, ionization products are measured directly; limited number of mechanisms for energy absorption
 - Results in measured variance **less** than that predicted by Poisson model
- Deviation from Poisson-predicted variance quantified in **Fano Factor, $F = \text{var}(N) / N$**
 - $F_{\text{semiconductor}} \sim 0.05 - 0.1$, $F_{\text{gas (ionization)}} \sim 0.2$
 - $F_{\text{scintillator}} \sim 1.0$



Fano Factor in Semiconductor Detectors

- [Spieler](#) provides concise, conceptual model for the origin of the Fano factor in semiconductor detectors
 - For more information, start with the [original paper from Fano](#)

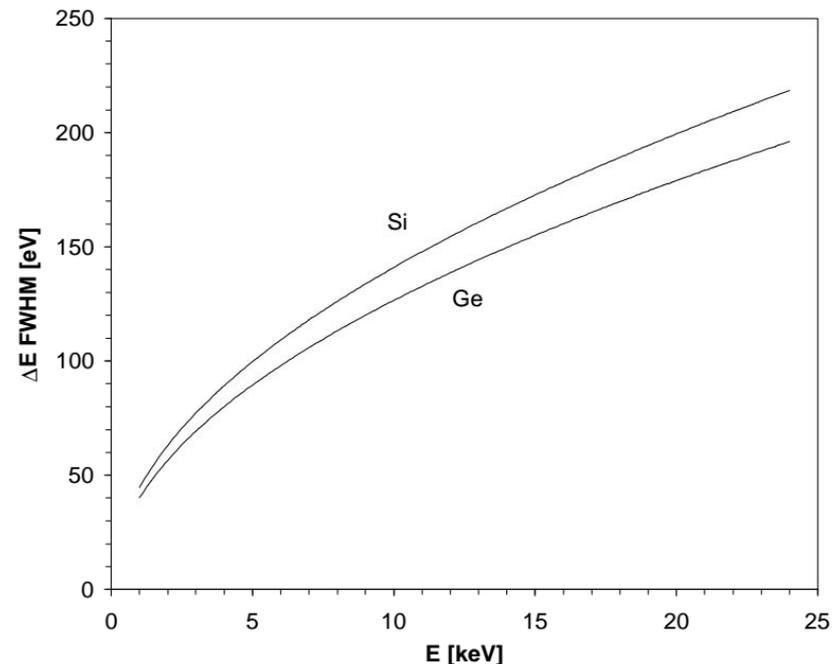
$$\Delta E = 2.35 \cdot \epsilon_i \sqrt{FN_Q} = 2.35 \cdot \epsilon_i \sqrt{F \frac{E}{w}} = \boxed{2.35 \cdot \sqrt{FE\epsilon_i}}$$

Si: $\epsilon_i = 3.6 \text{ eV}$ $F = 0.1$

Ge: $\epsilon_i = 2.9 \text{ eV}$ $F = 0.1$

- [Spieler 2.2.3](#)

Intrinsic Resolution of Si and Ge Detectors

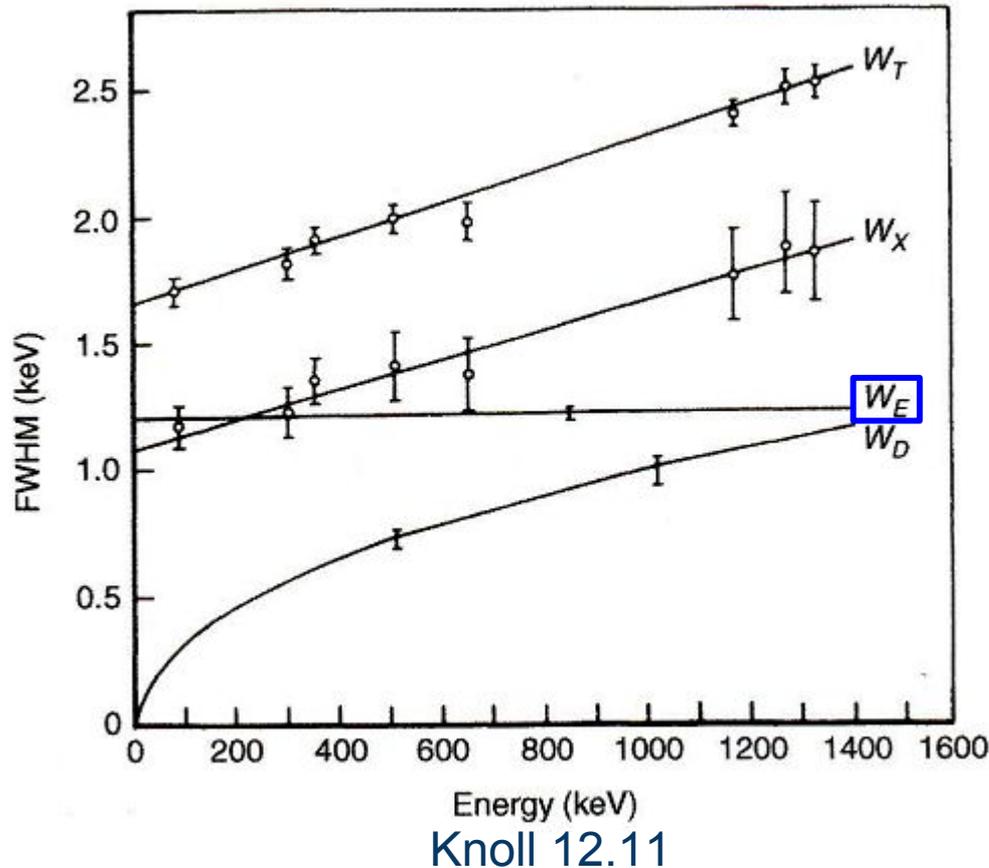


Energy Resolution - Statistics of Carrier Generation



$$(\Delta E_{\text{total}})^2 = (\Delta E_{\text{stat}})^2 + (\Delta E_{\text{electronics}})^2 + (\Delta E_{\text{charge-loss}})^2$$

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- Multiple elements contribute to width of peaks in energy spectrum
 - Statistics of carrier generation
 - **Electronic Noise**
 - Charge collection



Sources of Electronic Noise

- Basic mechanisms contributing to electronic noise (from [Spieler](#))

$$i = \frac{n e v}{l}$$

Labels for the equation:

- n : # Charges
- v : Carrier velocity
- l : Electrode coupling (Shockley-Ramo)
- i : Current (signal)

$$\langle di \rangle^2 = \frac{ne}{l} \langle dv \rangle^2 + \frac{ev}{l} \langle dn \rangle^2$$

- Current fluctuations given by fluctuations in **number of charge carriers** and **charge carrier velocity**
 - Velocity fluctuations** → thermal noise
 - Number fluctuations**
 - Shot noise, e.g. current flow through barrier junction
 - Carrier trapping/detrapping → 1/f noise



Quantifying Electronic Noise

- Describe noise in terms of spectral density, i.e. noise power per unit bandwidth

- Spectral noise power density
- Spectral noise voltage density
- Spectral noise current density

$$\frac{dP_n}{df}$$
$$\frac{dv_n^2}{df} = \frac{dP_n}{df} R$$
$$\frac{di_n^2}{df} = \frac{dP_n}{df} \frac{1}{R}$$

- Quantifying noise in terms of detector signal: **Equivalent Noise Charge**

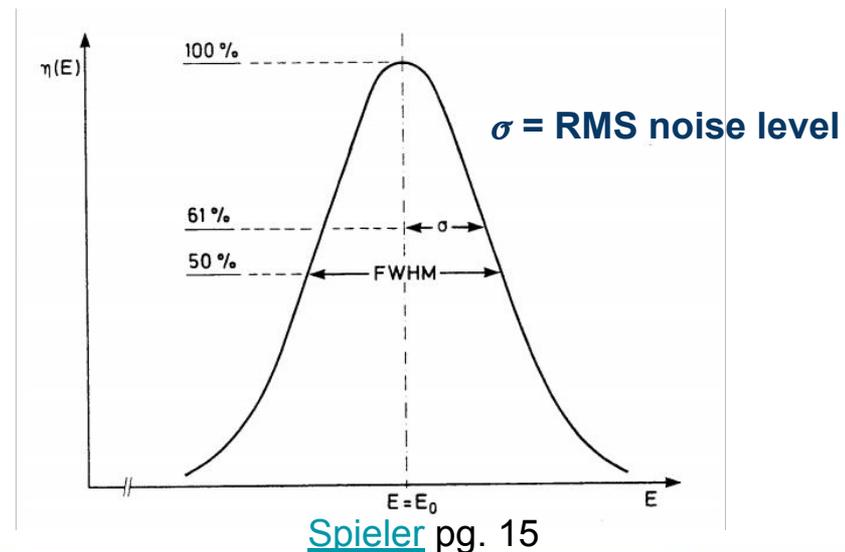
- Signal charge that yields $SNR = 1$
- For ionization detector with average ionization energy, ε_i , noise level can be expressed in terms of energy by

$$E_{noise}(eV) = \varepsilon_i \cdot ENC$$



Characteristics of Electronic Noise

- Spectral distribution of both **thermal** and **shot** is constant
 - I.e. “white” noise: $\frac{dP_{noise}}{df} = const.$
- 1/f noise exhibits a frequency dependence: $\frac{dP_{noise}}{df} = \frac{1}{f^a}$
 - $a \sim 0.5 - 2$ (discussed in a few slides)
- Thermal and shot noise are purely random and uncorrelated
- Amplitude distribution is Gaussian - deviations symmetric about DC baseline level





Thermal (Johnson) Noise

- Electron velocities given by thermal distribution
- Spectral density of noise power can be derived from long-wavelength approx. to blackbody spectrum (Spieler 3.4.1)

$$\frac{dP_n}{df} = 4kT$$

- In elements with finite resistance, gives rise to voltage/current fluctuations:

- Spectral noise voltage density: $\frac{dv_n^2}{df} \equiv e_n^2 = 4kTR$

- Spectral noise current density: $\frac{di_n^2}{df} \equiv i_n^2 = \frac{4kT}{R}$



Shot Noise

- Carrier injection over some potential barrier
 - Carrier injection across rectifying (PN junction), blocking contacts
 - Thermionic carrier generation (bulk leakage)
- Carrier injection an independent, stochastic process
 - Subject to statistical fluctuations
- Injection is random in time, contribution from each injection can be treated as delta pulse, yielding white (freq. independent) spectrum (see Spieler 3.4.2)

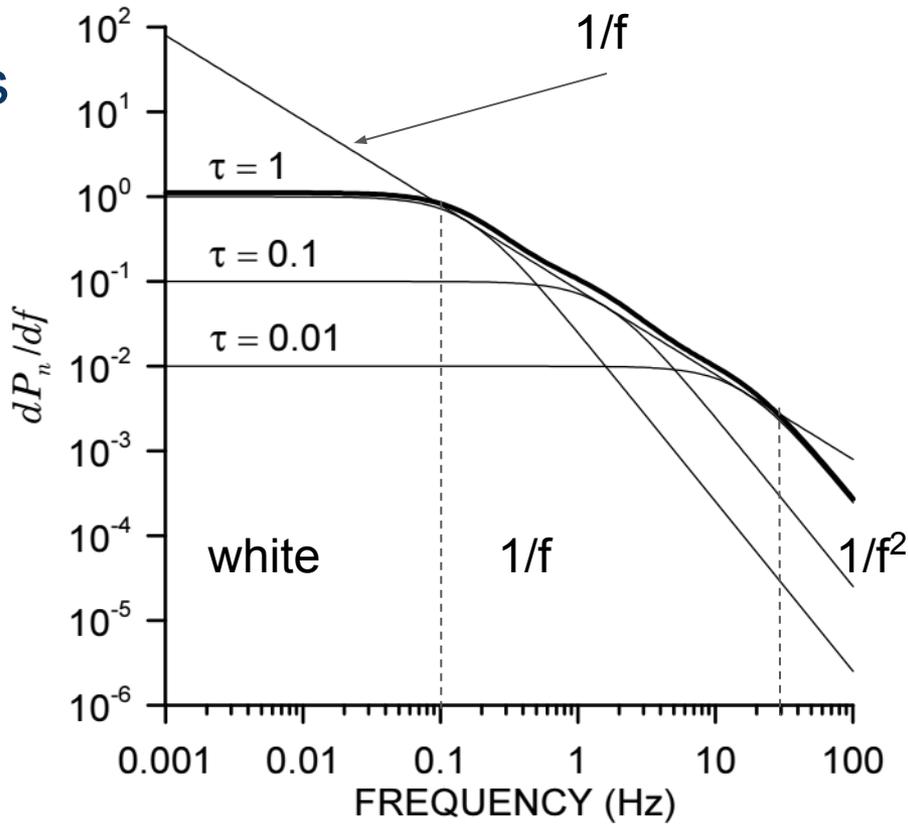
- Spectral current noise density:
$$i_n^2 \equiv \frac{di_n^2}{df} = 2eI$$

I = average current, Ne



1/f Noise

- Results from trapping/de-trapping of carriers
 - Trapping events are independent, random in time
- Characteristic times involved with traps of various depths
 - E.g. “shallow” trap → small τ , “deeper” trap → longer τ
- Multiple time constants give rise to 1/f behavior



$$i_{nf}^2 = 4NI^2 \left(\frac{\Delta G}{G} \right)^2 \frac{\tau}{1 + (\omega\tau)^2}$$

$$\frac{dP_{noise}}{df} = \frac{1}{f^a}$$

Where $a \sim 0.5 - 2$



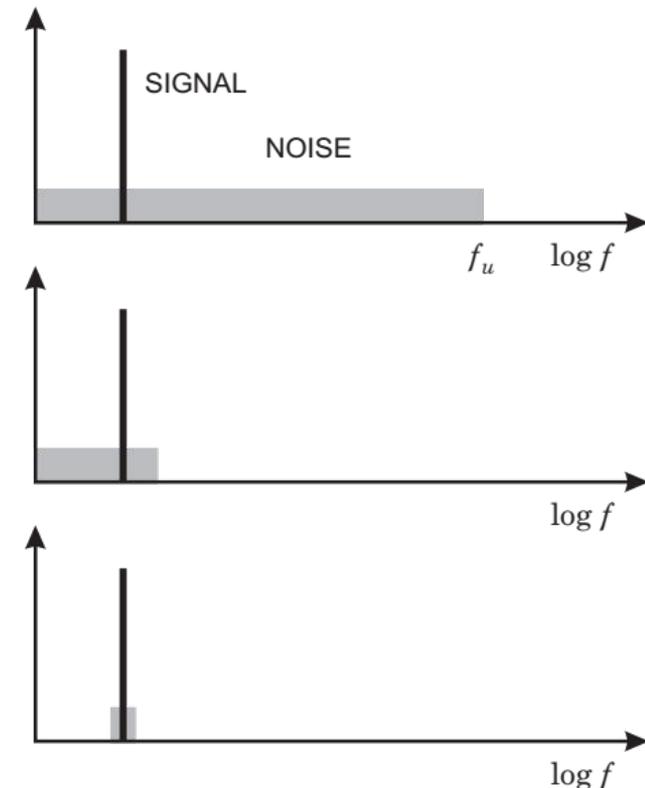
Signal Shaping

- Total noise of the system given by integrating spectral noise distribution over the bandwidth of the shaper

$$v_{no}^2 = \int_0^{\infty} e_n^2 A^2(f) df \quad \text{or} \quad i_{no}^2 = \int_0^{\infty} i_n^2 A^2(f) df$$

- v_{no}, i_{no} = noise at output of shaper
- $A(f)$ = gain of shaper

- Total noise increases with $(\text{BW})^{1/2}$
- **N.B.** Decreasing BW \rightarrow longer rise times

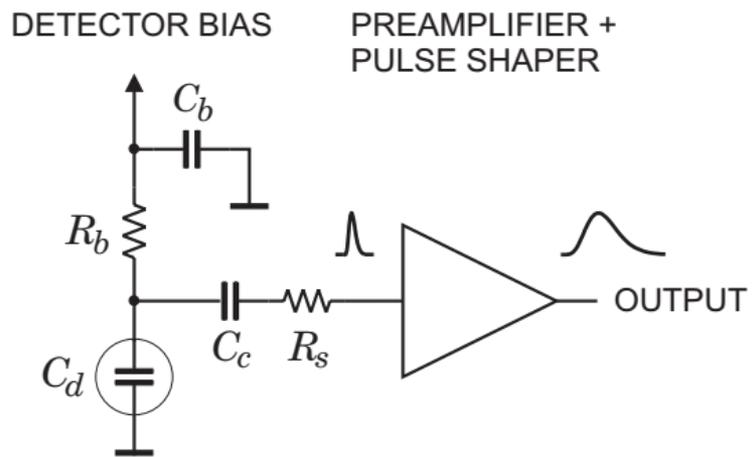


Spieler fig 3.3

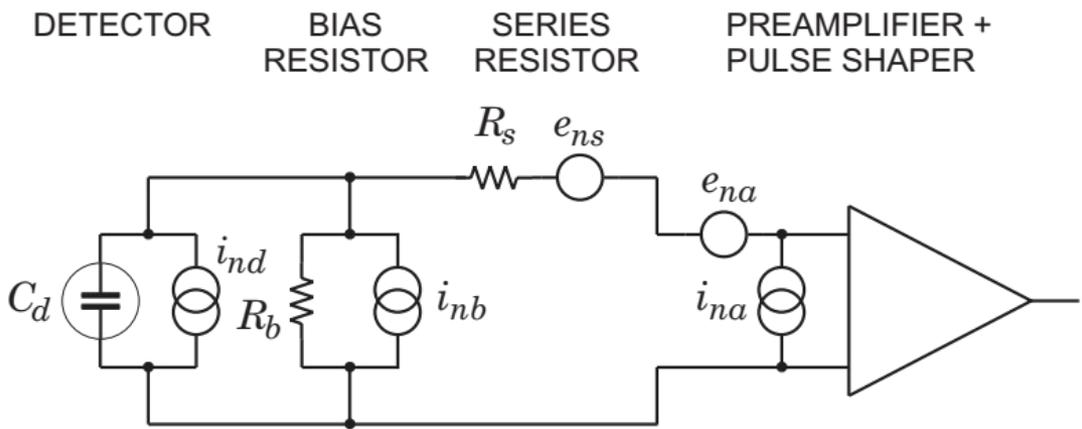


ENC of Charge-Sensitive Front End

Model Ckt of signal chain



Equiv. Ckt for analytic noise analysis



Spieler fig 4.7

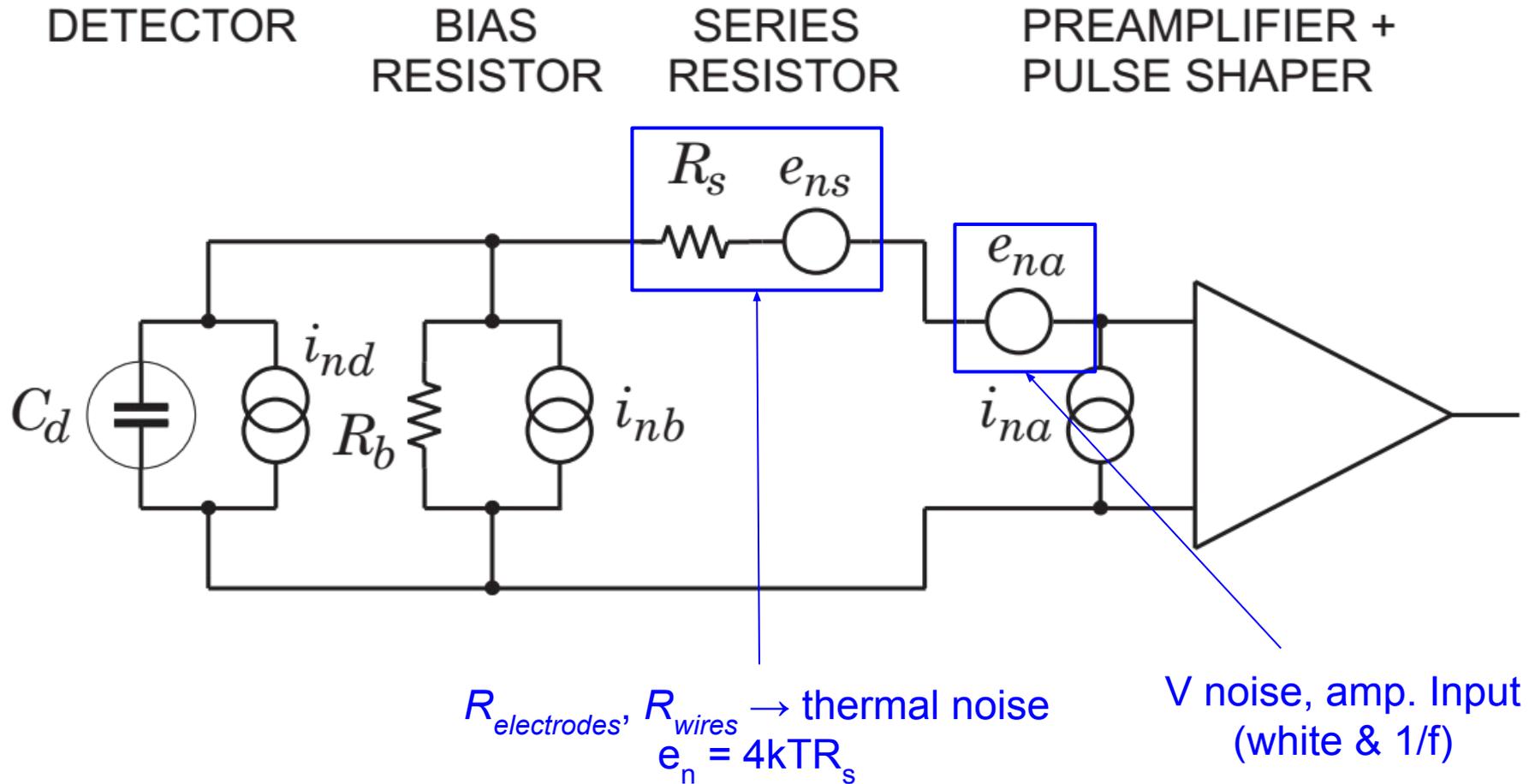
N.B. - Resistors can be modelled as voltage or current sources

- Resistors in **parallel** with input: noise current sources
- Resistors in **series** with input: noise voltage sources

This is where the “series” / “parallel” noise monikers come from



Voltage (Series) Noise Sources





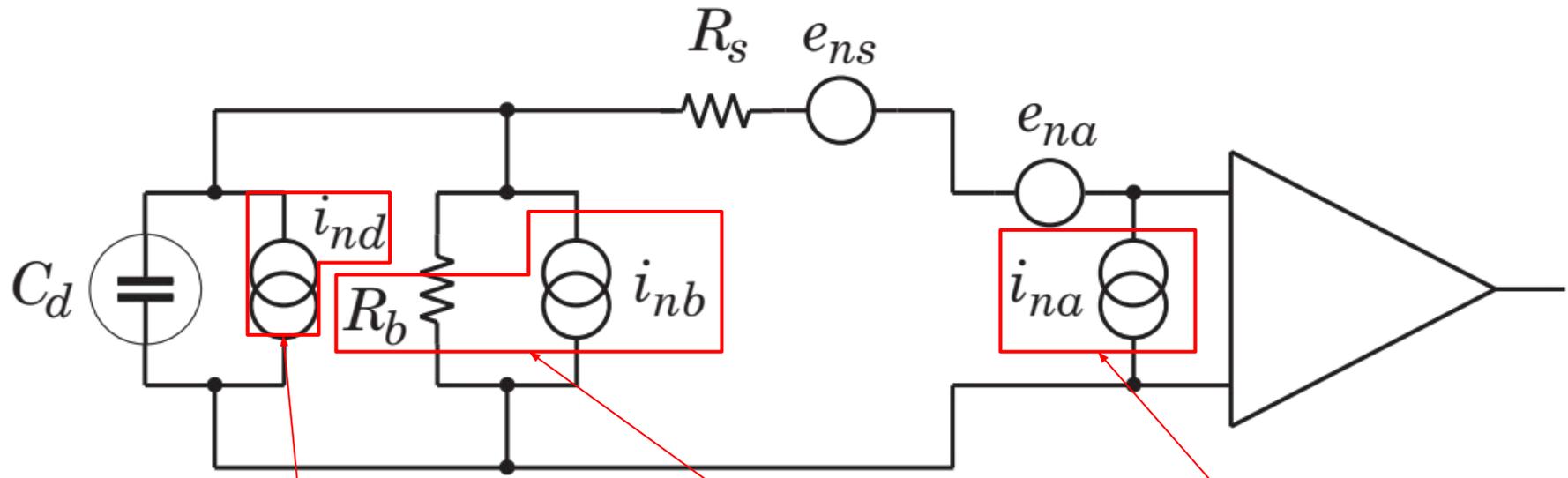
Current (Parallel) Noise Sources

DETECTOR

BIAS RESISTOR

SERIES RESISTOR

PREAMPLIFIER + PULSE SHAPER



Shot noise (e.g. leakage current)

$$e_{nd}^2 = i_{nd}^2 \frac{1}{(\omega C_d)^2} = 2eI_d \frac{1}{(\omega C_d)^2}$$

Thermal noise in

R_b

$$e_{np}^2 = \frac{4kT}{R_p} \left(\frac{R_p \cdot (-i/\omega C_d)}{R_p + (-i/\omega C_d)} \right)^2$$

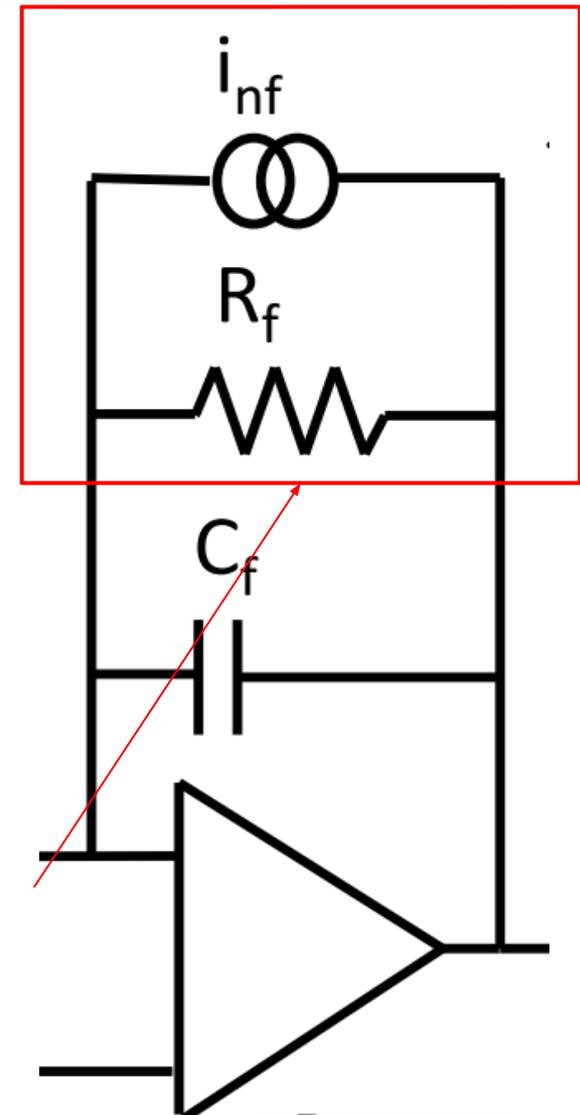
Shot noise in preamp input FET

$$\frac{i_{na}^2}{(\omega C_d)^2}$$



Current (Parallel) Noise Sources

- Noise from feedback resistor
 - Thermal (Johnson) noise in feedback resistor parallel to amplifier input
- Very low-noise front-ends try to reduce noise by moving away from resistive feedback to eliminate this noise source
 - Leakage current through diodes in TRPs are a noise source too!
 - Optical-reset CSA's



Thermal noise
 $i_n^2 = (4kT/R_f)$

Figure created by Brian Plimley



ENC Analysis

Cumulative input noise voltage

$$\begin{aligned} e_{ni}^2(f) &= e_{nd}^2 + e_{np}^2 + e_{nr}^2 + e_{na}^2 = \\ &= \frac{2eI_d}{(\omega C_d)^2} + \frac{4kTR_p}{1 + (\omega R_p C_d)^2} + 4kTR_s + e_{na} + \frac{i_{na}^2}{(\omega C_d)^2} \end{aligned}$$

Noise voltage at output depends on freq. resp. of amplifier

$$V_{no}^2 = \int_0^{\infty} e_{no}^2(f) df = \int_0^{\infty} e_{ni}^2(f) |A_v|^2 df$$

Specific example of **ENC for CR-RC** shaper with $\tau_{int} = \tau_{diff}$

$$Q_n^2 = \left(\frac{\epsilon^2}{8} \right) \left[\left(2eI_d + \frac{4kT}{R_p} + i_{na}^2 \right) \cdot \tau + (4kTR_s + e_{na}^2) \cdot \frac{C^2}{\tau} + 4A_f C^2 \right]$$

See Spieler 4.3.5 for full treatment

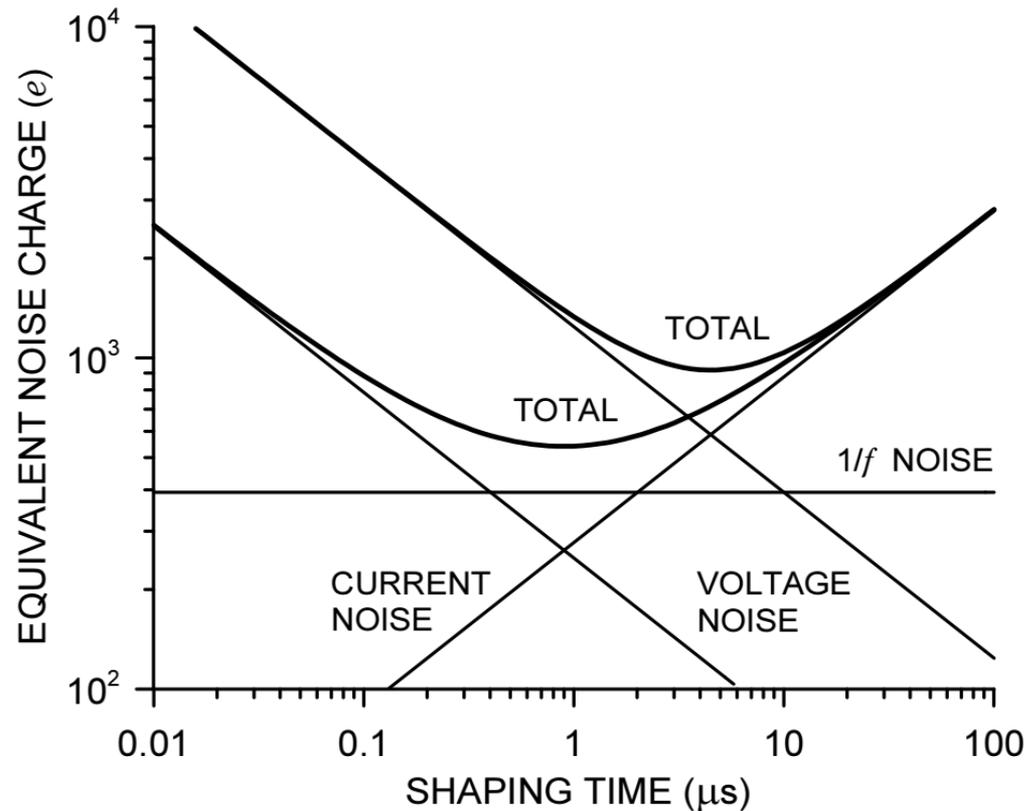


ENC Analysis

A more general result for ENC was derived in RA 4 (Radeka):

$$Q_n^2 = i_n^2 F_i T_S + e_n^2 F_v \frac{C^2}{T_S} + F_{vf} A_f C^2$$

- T_s = shaping time
- C = total input capacitance
- F_i, F_v, F_{vf} = “Shape factors” - can be computed from IR of shaper (see Spieler 4.5.2)



Spieler fig 4.29



ENC Analysis

A more general result for ENC analysis was derived in reading assignment 4:

$$Q_n^2 = \boxed{i_n^2 F_i T_S} + \boxed{e_n^2 F_v \frac{C^2}{T_S}} + \boxed{F_{vf} A_f C^2}$$

Current noise sources

- Independent of input capacitance
- Contribution increases with increasing T_s

Voltage noise sources

- Increases rapidly with input capacitance
- Contribution decreases with increasing T_s

1/f noise sources

- Increases rapidly with input capacitance
- Independent of T_s , but depends on BW

Minimum noise:

$$Q_n^2 = 2e_n i_n C \sqrt{F_i F_v} + F_{vf} A_f C^2$$

$$v_n^2 = \int_{f_l}^{f_u} \frac{A_f}{f} df = A_f \log \frac{f_u}{f_l}$$

Note dependence on shape factors!



Noise Curves from Real Systems

- Some examples in Spieler 4.4
- Noise curve from one of our HPGe's - can you diagnose the problem?

Shaping Time	FWHM (keV) @ 2402.57 keV, HV = +2400	FWHM @ 2402 HV = +2500
12	0.75	4.6
6	0.69	2.6
4	0.74	2.05
2	0.85	1.33
1	1.01	1.11
0.5	1.27	1.27

